05.1—73   POLYMORPHIC TRANSITIONS IN ORGANIC CRYSTALS.
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Experimental study of the mechanism of temperature poly-
morphic transitions (P.T.) in organic crystals involves
several stages: 1) obtaining perfect single crystal sam-
ples; 2) study of the phase behaviour and morphology of
growth of a new phase; 3) microcalorimetric study;
4) determination of orientational phase interrelations;
5) structural interface studies; 6) estimation of thermo-
dynamic parameters; 7) modelling of the mechanism of
P.T. Studies of the polymorphic transitions carried out
along these lines on a number of substances indicate
that P.T. belong to the first order phase transitions
and represent a growth of crystals of   a new phase in-
side the matrix phase. The defects of the initial crys-
tal are the centers of nucleation. All P.T. are charac-
terized by hysteresis, the actual value of which depends
on macroscopic parameters as well as on specificity of
the structure of the substance under study. Orientation
of crystals of a new phase inside the matrix phase may
be different depending mainly on the geometric form of
the molecule. Specific intermolecular forces (H-bonds),
conformational mobility of molecules etc. complicate
the picture of the process and the possibility of modell-
ing the mechanism of P.T. A number of phenomena accompa-
nying P.T. are investigated, their emergence is explain-
ed. Microcalorimetric studies and calculations of cha-
racteristic temperatures describe, to a rather high
accuracy, the type of polymorphic transitions and help
in understanding the general idea of the phenomenon. The
revealed factors allow semiempirical prediction of some
main features of P.T. in majority of molecular crystals.
To elucidate whether there is a direct relation between
morphological, structural and energetic characteristics
on the one hand and molecular mechanism of the phenome-
non on the other, one should have at least by one order
of magnitude more statistic data on P.T. than available.
The significance of the revealed factors for modelling
structural transitions in heteromolecular complexes is
discussed.

05.1—74   STRUCTURAL INTERPRETATION OF PHASE TRANSI-
TIONS IN SODIUM CARBONATE. By P.M. de Wolff and
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The Netherlands.

The two second-order transitions α-β (489°C) and β-γ
(361°C) of sodium carbonate can be interpreted in
terms of a columnar structure model. This view is
based on detailed structure analyses at various
temperatures and it is supported by analogous phase
changes in rubidium- and potassium carbonate. In the
three isomorphous α phases, space group P6$_3$mmc, face-
sharing A'O$_6$ octahedra form parallel columns in a
hexagonal array, with C and the remaining alkali atoms
A'' alternating in a row on each axis between three
adjoining columns. So the latter are connected by the
anions through oxygen sharing; in the α form the
anions are perpendicular to the columns. The oxygens
act as flexible joints allowing huge shear distortions
in the sheaf of undistorted columns. This is what
happens when thermal motion of the weakly bound A''
ions becomes insufficient to stabilize the α form. For
Na$_2$CO$_3$ a the α-β softening mode is one component of E$_{1g}$
(rotation of anions about an O-O edge); it couples at
$q \approx 0$ with the a-polarized shear belonging to E$_{1u}$
(axes: b $\approx$ a$\sqrt{3}$). In the γ phase the complementary
modes of E$_{1g}$ and E$_{1u}$ (wiggle of anions about C-O and
b-polarized periodic shear) are condensed with an
incommensurate q. The columnar model allows an
interpretation of the anomalous thermal expansion and
of the fact that q is normal to the average plane of
the anions. Diffuse scattering shows persistance of
the softening throughout the β-phase. The thermal

behaviour of birefringence and IR spectra is readily
interpreted in terms of the known structural changes.

05.1—75   LATTICE PARAMETERS AND THERMAL EXPANSION OF
V$_3$O$_5$ NEAR PHASE TRANSITION. By S. Åsbrink* and
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holm, Sweden* and Institute for Low Temperature and
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Poland**.

V$_3$O$_5$ undergoes a first order phase transition at T$_t$ =
427.9 K of the semiconductor-to-metal type. Both low-
and high-temperature phases are monoclinic with Z = 4
and the space groups P2/c and I2/c, respectively. Ac-
cording to the accurate structure determination
(Åsbrink, Acta Cryst. (1980) B36,1332; Hong & Åsbrink,
Acta Cryst. (1982) B38,713), an ordered distribution of
tri- and tetravalent vanadium atoms in low-V$_3$O$_5$ is re-
placed by the only partially ordered one in high-V$_3$O$_5$,
but the shifts in the atomic positions are less than
0.1 Å.
A dependence of the lattice parameters vs. temperature
was precisely investigated using Bond-type diffracto-
meter with high-temperature attachment (Łukaszewicz et
al., Krist. Tech. (1978) 13,561). The measurements were
performed in the temperature range 298-480 K with the
stability better than ± 0.02 K. The results for the
temperatures just below and just above T$_t$ are following:

	298.16 K	427.85 K	427.90 K
a (Å)	9.85727(4)	9.85623(4)	9.84816(4)
b (Å)	5.04057(2)	5.04135(2)	5.03457(2)
c (Å)	6.98882(5)	6.99646(5)	7.00225(5)
β (°)	109.4742(6)	109.4815(6)	109.4947(6)
V (Å³)	327.383(6)	327.742(6)	327.278(6)

The linear Lagrangian strain tensor at the transition
was calculated and diagonalized (Schlenker et al.,
Acta Cryst. (1978) A34,52). Its principal components
are following:
$$l_1 = -9.18 \cdot 10^{-4}, \quad l_2 = -13.42 \cdot 10^{-4}, \quad l_3 = 8.45 \cdot 10^{-4}$$
$\upsilon = 5.77°$ (axes description - see scheme in fig.)
The thermal expansion tensor was applied in the limit
of infinitesimal displacements. Its coefficients α_i and
angle υ are plotted below as the functions of tempera-
ture.
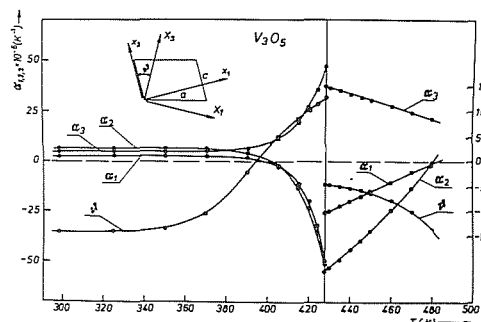
One can conclude that the crystal structure starts to
change continuously at the temperature of about 360 K.
Above T$_t$ in the range of 50 K the thermal parameters
tend to return to their previous values from the low-
temperature phase. The analysis of the strain as well
as the thermal expansion tensor leads to the conclusion
that the directions of the structure instabilities lie
closely to the direction parallel to the c-axis (i.e.
along the diagonal of the basis of the VO$_6$ octahedra)
and to the direction perpendicular to it. Referring to
the structural results one may suggest that from 360 K
upwards the difference between the two unequivalent va-
nadium positions regarding the valencies starts to dis-
appear. At T$_t$ the valency-mixing reaches a level where
a symmetry change is needed for the energetical reasons
and with further temperature increase results in the
statistical distribution of the vanadium atoms of dif-
ferent valencies.