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## 12-Amorphous, Imperfectly Ordered and Quasi-periodic Materials

**PS-12.02.13 DISORDER IN THE PLASTIC PHASE OF HEXACHLOROETHANE.** By A.Criado and A.Muñoz, Departamento de Física de la Materia Condensada, Instituto de Ciencia de Materiales, Universidad de Sevilla, Spain.

Hexachloroethane is a molecular crystal which shows an orientationally disordered cubic phase between 344K and the melting point at 458K. Neutron diffraction experiments (Gerlach, Hohlwein, Prandl & Schulz, 1981, Acta Cryst. A37, 904- 908) show that the orientational probability is maximum along the [100] and [111] direction for the chlorine and carbon atoms respectively. X-ray diffuse scattering (Gerlach & Prandl, 1988, Acta Cryst. A44, 128- 135) and quasielastic neutron scattering (Gerlach, Dorner, Prandl & Lefebvre, 1988, Acta Cryst. A44, 251- 257) work has also been carried out.

In this work we present a molecular dynamics simulation of the plastic phase using intermolecular potentials in order to explain the nature of the disorder. The calculated atomic probability distributions are in good agreement with the experiment. The molecular probability distribution, not accessible by diffraction methods, is also obtained. The origin of the disorder is studied from the correlations of the molecules with their surrounding neighbours, which reflect in the diffuse scattering patterns.

### PS-12.02.14 MODEL OF THE SHORT-RANGE ORDER OF AMORPHOUS IRON AND ITS ENERGETIC STABILITY.

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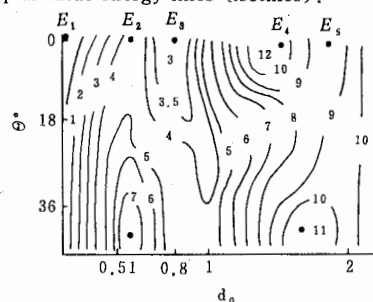
Diffraction pattern from iron amorphized by boron is recorded using X-ray diffractometry. The Interference component was separated after different corrections dependent on radiation and substance properties, and normalization of this pattern has to be performed. Fourier transform of this component gave the atomic radial distribution function based on which the short-range order model was obtained, namely, icosahedron with the nearest interatomic distance between Fe atoms equal to 2.52 Å. Interference components of amorphous Fe were noticeably different before and after annealing suggesting the displacement of the upper part of icosahedron relative to the lower one. Icosahedral short-range order was confirmed by calculation of the theoretical interference component, such a displacement giving markedly different results. In this connection, fit of the icosahedron model to amorphous Fe was considered in the energetic context.

Configuration (potential) energy of cluster - icosahedron was calculated on the basis of coordinates with Lennard-Jones type atomic interaction (the interaction type was chosen the structure of the short-range order to be due to interaction energy (potential pit) at the nearest interatomic distances). The energy of the system was calculated in approximation of atomic pair additivity to L-J potential (6-12):

$$E_N = \sum_{i \neq j} \varphi(r_{ij}), \quad \varphi(r) = \epsilon \left[ \left( \frac{r_1}{r} \right)^{12} - 2 \left( \frac{r_1}{r} \right)^6 \right],$$

where  $\epsilon$  - the deep of potential pit,  $r_1$  - the atomic diameter taken

as a length unit. Icosahedron centered at the origin was arranged so that two opposite tops were located at vertical axis and the plane from which polar angle  $\theta$  was measured has gone through one of the equatorial atoms in the upper half of icosahedron. The six upper atoms were displaced by  $d$  relative to the lower ones. Energy values of icosahedra deformed were calculated as functions of values of displacement  $d$  and an angle  $\theta$  of turn relative to an axis. The potential energy surface represented the picture of equal value energy lines (isolines):



On this surface there was a stable local minimum at  $E_0 = -2.52$ /atom except global minimum at  $E_0 = -2.79$ /atom and the shortest distance between them has gone through the saddle point  $E_1 = -2.39$ /atom. Consequently, in the case of Fe the potential barrier value is equal to  $0.13\epsilon \approx 340$  K at  $\epsilon = 2600$  K. Notice that points  $E_0$  and  $E_1$  at the energy surface correspond to the shortrange orders in icosahedron and icosahedron with displacement.

### PS-12.02.15 TOPOLOGICAL EQUIVALENCE CLASSES OF A TWO DIMENSIONAL RANDOM NETWORK WITH COORDINATION NUMBER 3 ON THE BASE OF FIBRE BUNDLES.

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The translational invariance principle is known to be not adequate for the description of disordered solid's structure. Therefore we propose a new approach based on the local invariance principle which has an empirical ground - the existence of short range order (SRO). So the local approach is available not only for the description of amorphous solids but crystalline [1] and quasicrystalline [2] ones too. It turns out that the local approach can be formulated in natural way within a mathematical tool of a fibre bundle theory.

As an example we consider a simplest case, namely the model of a two dimensional random network with coordination number three. To construct the corresponding fibre bundle we introduce the notion of the short range order configuration (SROC), which differs from that of the SRO essentially. The set of all of the SROC's which can be realized in the structure forms the short range order configuration space (SROCS). We have shown that the principal bundle  $P$  and the associated one  $E$  have the form:

$$P(S^2, SO_0(1,2) \times SO(2)), \\ E(S^2, D^2 \times S^1, SO_0(1,2) \times SO(2))$$

respectively. The basis  $S^2$  (two dimensional sphere) arises from boundary conditions which must be imposed on an atomic system. In this respect we used the isotropy property of most glasses. The fibre (the SROCS)  $D^2 \times S^1$  is a direct product of the